

BOUNDARY VALUE PROBLEM IN LOWER SEMIFIELDS DEVIATING FROM CHARACTERISTICS FOR A PARABOLIC- HYPERBOLIC EQUATION

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Annotation: This article proves the existence and uniqueness of a solution to a nonlocal conditional boundary value problem for a boundary value problem in lower semifields containing lines shifted from the characteristic for an equation of parabolic-hyperbolic type.

Key words: Equation of parabolic-hyperbolic type, characteristic triangle, usual solution, integral energy method, simple solution, Green's function, Volterra integral equation of the second kind, D'Alembert's formula.

1. Setting the issue

Consider the following equation:

$$0 = Lu \equiv \begin{cases} u_{xx} - u_y, & (x, y) \in \Omega_0, \\ u_{xx} - u_{yy}, & (x, y) \in \Omega_j (j = \overline{1, 3}), \end{cases} \quad (1)$$

here Ω_0 as a field $x > 0, y > 0$ when $y = 0, x = 1, y = 1, x = 0$ located in straight lines AB, BB_0, B_0A_0, A_0A , a rectangular area bounded by sections, Ω_1 field $x < 0, y > 0$ da ΔAA_0D is located inside the characteristic triangle $AK: x = \gamma_1(y)$ smooth curve and equation (1). $BP: y - x = 1$ area limited by characteristic, Ω_2 field in $x > 0, y > 0$ ΔBB_0E is located inside the characteristic triangle $AC: x = -\gamma_2(y)$ smooth curve and equation (1). $B_0M: x + y = 2$ area limited by characteristic, Ω_3 field in $x < 0, y < 0$ (1) of the equation $AC: y + x = 0$ va $BC: y - x = -1$ limited by the characteristic ΔAA_0D characteristic triangular area,

Let us introduce the following definitions: $J_1 = \{(x, y): 0 < x < 1, y = 0\}$,

$J_2 = \{(x, y): x = 0, 0 < y < 1\}, J_3 = \{(x, y): x = 1, 0 < y < 1\}$,



$$\Omega = \Omega_0 \cup \Omega_1 \cup \Omega_2 \cup \Omega_3 \cup J_1 \cup J_2 \cup J_3, C\left(-\frac{1}{2}, \frac{1}{2}\right), D\left(-\frac{1}{2}, \frac{1}{2}\right), E\left(\frac{3}{2}, \frac{1}{2}\right),$$

$$\theta\left(\frac{x}{2}; -\frac{x}{2}\right), \left[\theta^*\left(\frac{\lambda(x)+x}{2}; \frac{\lambda(x)-x}{2}\right)\right]$$

there $\theta_2(x) \left[\theta_2^*(x)\right]$ (1) of Eq $(x, 0) \in J_3$ with point-to-point characteristics $AC[AN]$ the coordinates of the intersection point of the characteristics.

A-Condition. $y = -\gamma_1(x)$ va $\gamma_j(y) (j = 2, 3)$ - let the given functions fulfill the following conditions:

1) $\gamma_1(x)$ va $\gamma_j(y) (j = 2, 3)$ functions respectively ΔACB and $\Delta AA_0D, \Delta BEB_0$ should be completely inside the characteristic triangles;

2) $\gamma_1(x) \in C^2(0, 1), \gamma_j(y) \in C^2(0, 1) (j = 2, 3)$ be relevant;

3) $t \pm \gamma_j(t) (j = \overline{1, 3})$ - monotonically growing;

4) $\gamma_1(0) = 0, \gamma_2(0) = -1, l_1 + \gamma_1(l_1) = 1, l_2 - \gamma_2(l_2) = 2, l_3 + \gamma_3(l_3) = 1, l_j = const$

$$l_j \in \left(\frac{1}{2}; 1\right).$$

Description. (1) as a regular solution of Eq Ω ,

$$W_1 = \left\{u : u(x, y) \in C^1(\overline{\Omega}) \cap C_{x,y}^{2,1}(\Omega_0) \cap C^2(\Omega_i), i = \overline{1, 3}\right\}$$

belonging to the class and $\Omega_i (i = \overline{0, 3})$ in the field (1) satisfying Eq $u(x, y)$ the solution is told.

I – Matter. Find a regular solution of equation (1) satisfying the following conditions:

$$\left[u_x - u_y\right] \theta(x) + \mu(x) \left[u_x - u_y\right] \theta^*(x) = \varphi(x); \quad (2)$$

$$u|_{A_0D} = g_1(y); u|_{B_0E} = g_2(y) \quad (3)$$

$$\left[u_x + u_y\right] \Big|_{AD} = p(y); \quad (4)$$

$$\left[u_x - u_y\right] \Big|_{BE} = q(y); \quad (5)$$

$$u(A) = u(B) = 0; \quad (6)$$

There $\mu(x), \varphi(x), g(y), p(y)$ va $q(y)$ - are sufficiently smooth functions given.

Theorem. If $\mu(x) \neq -1; \mu(x), \varphi(x) \in C^1[0, 1] (i = \overline{1, 3}), g(y), p(y), q(y) \in C^2(0; 1)$

and **A** the conditions are met, **I**-is the only regular of the problem



there will be a solution.

Proof. Ω_1 The appearance of the solution of the Cauchy problem in fields is as follows.

$$u(x, y) = \frac{1}{2} \left[\tau_1(x+y) + \tau_1(x-y) + \int_{x-y}^{x+y} v_1(t) dt \right], \quad 0 < x < 1$$

From this $\theta\left(\frac{x}{2}; -\frac{x}{2}\right)$ and $\theta^*\left(\frac{\lambda(x)+x}{2}; \frac{\lambda(x)-x}{2}\right)$ we get the following functional relations

at the points:

$$(1 + \mu(x))v_1(x) = [1 + \mu(x)]\tau_1'(x) - \varphi(x), \quad 0 < x < 1 \quad (7)$$

From (1) Ω_0 in the field $y \rightarrow +0$ passing to the limit, we get the following equation.

$$(\tau_1(x))'' = v_1(x) \quad (8)$$

Substituting (7) into (8). We can form a second-order ordinary differential equation with respect to $\tau_1(x)$.

$$(\tau_1(x))'' - \tau_1'(x) = -\frac{\varphi(x)}{1 + \mu(x)} \quad (9)$$

(9) equation $\tau_1(0) = 0$ and $\tau_1(1) = 0$ the shape of the general appearance of the picture under the conditions.

$$\tau_1(x) = \int_0^x \frac{\varphi(t)[1 - e^{x-t}]}{1 + \mu(t)} dt + \frac{e^x - 1}{e - 1} \int_0^1 \frac{\varphi(t)[e^{1-t} - 1]}{1 + \mu(t)} dt \quad (10)$$

using (8), we find $v_1(x)$

$$v_1(x) = -\int_0^x \frac{\varphi(t)e^{x-t}}{1 + \mu(t)} dt + \frac{e^x}{e - 1} \int_0^1 \frac{\varphi(t)(e^{1-t} - 1)}{1 + \mu(t)} dt - \frac{\varphi(x)}{1 + \mu(x)} \quad (11)$$

We prove the existence of a solution to the problem by the method of integral equations.

For this, we use the functional relations (8) - (9) and the solution of the first boundary value problem for the equation (1) in domain Ω_0

$$u(x, y) = \int_0^1 \tau_1(t)G(x, y; t, 0)dt + \int_0^y \tau_2(t)G_t(x, y; 0, z)dz - \int_0^y \tau_3(z)G_t(0, y; 1, z)dz, \quad (12)$$

it will be visible.



There $G(x, y; t, z) = \frac{1}{2\sqrt{\pi(y-z)}} \sum_{n=-\infty}^{\infty} \left[e^{-\frac{(x-t+2n)^2}{4(y-z)}} - e^{-\frac{(x+t+2n)^2}{4(y-z)}} \right]$ - Green's function of the first

boundary value problem for the heat transfer equation.

$\tau_k(y), v_k(y)$ ($k=2,3$) to get the relation between the functions, differentiate once: We make

$$u_x(x, y) = \int_0^1 \tau_1(t) G_x(x, y; t, 0) dt + \int_0^y \tau_2(z) G_{tx}(x, y; 0, z) dz - \int_0^y \tau_3(z) G_{tx}(x, y; 1, z) dz, \quad (13)$$

The following $N(x, y; t, z) = \frac{1}{2\sqrt{\pi(y-z)}} \sum_{n=-\infty}^{\infty} \left[e^{-\frac{(x-t+2n)^2}{4(y-z)}} + e^{-\frac{(x+t+2n)^2}{4(y-z)}} \right]$ if we define,

$G_{tx}(x, y; t, z) = N_z(x, y; t, z), G_x(x, y; t, z) = -N_t(x, y; t, z)$ We will have relationships.

According to definition, $u_x(0, y) = v_2(y)$ we get the following relations from (13):

$$v_2(y) = \int_0^1 \tau_3'(z) N(x, y; 1, z) dz - \int_0^y \tau_1'(t) N(x, y; t, 0) dt + \int_0^y \tau_2'(z) N(0, y; 0, z) dz$$

$$v_3(y) = \int_0^1 \tau_3'(z) N(1, y; 1, z) dz - \int_0^y \tau_1'(z) N(1, y; t, 0) dz + \int_0^y \tau_2'(z) N(1, y; 0, z) dz \quad (14)$$

It is known, The general solution of equation $u_{xx} - u_{yy} = 0$ is in the form

$$u(x, y) = f_1(x+y) + f_2(x-y) \quad (15)$$

in this $f_1(\cdot), f_2(\cdot)$ - second-order continuous differentiation is an unknown function.

From condition (4) and (15), we get $f_1'(y - \gamma_2(y)) = p(y), 0 \leq y \leq l$, from Eq $y - \gamma_2(y) = t$

solve looking for $y = \delta_1(t)$

$$f_1'(t) = \frac{1}{2} p(\delta_1(t)), 0 \leq y \leq l, \text{ From this}$$

$$f_1(y) = f_1(0) + \frac{1}{2} \int_0^y p(\delta(t)) dt, 0 \leq y \leq l.$$

From condition (3.8) and from (15). we will have $f_2'(y - \gamma_2(y)) = q(y), 0 \leq y \leq l$, from Eq

solve $y - \gamma_2(y) = t$ looking for $y = \delta_1(t)$

$$f_2'(t) = \frac{1}{2} q(\delta_1(t)), 0 \leq y \leq l, \text{ From this}$$



$$f_2(y) = f_2(0) + \frac{1}{2} \int_0^y q(\delta(t)) dt, \quad 0 \leq y \leq l.$$

Now in $l \leq y \leq 1$ $u|_{A_0D} = g_1(y)$ and given $u|_{B_0E} = g_2(y)$ conditions

$$\begin{cases} f_1(y) = g_1\left(\frac{y-1}{2}\right) + f_2(1), & l \leq y \leq 1 \\ f_2(y) = g_2\left(\frac{2-y}{2}\right) + f_1(2), & l \leq y \leq 1 \end{cases}$$

To (15) we put the value of $f_1(y)$ and $f_2(y)$ and get the following.

$$u(x, y) = \begin{cases} f_2(x-y) + \frac{1}{2} \int_0^y p(\delta(t)) dt + f_1(0), & 0 \leq y \leq l, \\ f_2(x-y) + g_1\left(\frac{y-1}{2}\right) + f_2(1), & l \leq y \leq 1. \\ f_1(x+y) + \frac{1}{2} \int_0^y q(\delta(t)) dt + f_2(0), & 0 \leq y \leq l, \\ f_1(x+y) + g_2\left(\frac{2-y}{2}\right) - f_1(2), & l \leq y \leq 1. \end{cases} \quad (16)$$

$u_y(0, y) = \tau'_i(y)$, $i = 2, 3$, If we differentiate the equality (16) once with respect to y , we say

$x \rightarrow 0$

$$\tau'_2(y) = \begin{cases} -f'_2(-y) + \frac{1}{2} p(\delta(y)), & 0 \leq y \leq l, \\ -f'_2(-y) + \frac{1}{2} g'_1\left(\frac{y-1}{2}\right), & l \leq y \leq 1. \end{cases} \quad (17)$$

And at $x \rightarrow 1$

$$\tau'_3(y) = \begin{cases} f'_1(1+y) + \frac{1}{2} q(\delta(1+y)), & 0 \leq y \leq l, \\ f'_1(1+y) - \frac{1}{2} g'_2\left(\frac{1-y}{2}\right), & l \leq y \leq 1. \end{cases} \quad (18)$$

Putting (16) into (17) and (18). $f'_i(y)$ $i = 2, 3$ for the function $\tau_i(y)$ $i = 2, 3$ and $v_i(y)$ $i = 2, 3$ we get the following functional relation between the functions:



$$\begin{cases} \tau_2'(y) = v_2(y) + p(\delta(y)), & 0 \leq y \leq l, x < 0 \\ \tau_2'(y) = v_2(y) + g_1\left(\frac{y-1}{2}\right), & l \leq y \leq 1, x < 0 \end{cases} \quad (19)$$

$$\begin{cases} \tau_3'(y) = v_3(y) + q(\delta(1-y)), & 0 \leq y \leq l, x > 1 \\ \tau_3'(y) = v_3(y) + g_2\left(\frac{1-y}{2}\right), & l \leq y \leq 1, x > 1 \end{cases} \quad (20)$$

From (19) and (20). $v_2(y)$ and If we put $v_3(y)$ in (14)

$$\begin{aligned} \tau_2'(y) - \int_0^1 \tau_3'(z) N(x, y; 1, z) dz + \int_0^y \tau_1'(t) N(x, y; t, 0) dt - \int_0^y \tau_2'(z) N(0, y; 0, z) dz &= F_2(y) \\ \tau_3'(y) - \int_0^1 \tau_3'(z) N(1, y; 1, z) dz + \int_0^y \tau_1'(z) N(1, y; t, 0) dz - \int_0^y \tau_2'(z) N(1, y; 0, z) dz &= F_3(y) \end{aligned} \quad (21)$$

There

$$\begin{aligned} F_2(y) &= \frac{1}{2} p(\delta(1-y)) - \frac{1}{2} g_1\left(\frac{1-y}{2}\right) \\ F_3(y) &= \frac{1}{2} q(\delta(1-y)) - \frac{1}{2} g_2\left(\frac{1-y}{2}\right) \end{aligned}$$

We get the system (21) in the form

$$\begin{cases} \tau_2'(y) - \int_0^y \tau_2'(z) N(0, y; 0, z) dz = F_2^*(y), \\ \tau_3'(y) + \int_0^y \tau_3'(z) N(1, y; 1, z) dz = F_3^*(y). \end{cases} \quad (22)$$

Where

$$\begin{cases} F_2^*(y) = F_2(y) + \int_0^y \tau_3'(z) N(x, y; 1, z) dz - \int_0^1 \tau_1'(t) N(x, y; t, 0) dt, \\ F_3^*(y) = F_3(y) + \int_0^y \tau_2'(z) N(1, y; 0, z) dz - \int_0^y \tau_1'(z) N(1, y; t, 0) dt \end{cases} \quad (23)$$

$$(14) \text{ system } |N(0, y, 0, z)| \leq \frac{2}{\sqrt{\pi|y-y_1|}} \sum_{n=1}^{\infty} \left| e^{-\frac{n^2}{|y-y_1|}} \right| \leq \text{const}$$

$|F_2^*(y)| \leq \text{const}$ because we solve equation 1 from the system (23) by the method of successive approximation and get

$$\tau_2(y) = F_2^*(y) + \int_0^y F_2^*(t) K(t, y) dt \quad (24)$$



putting $F_2^*(y)$ in (24).

$$\begin{aligned} \tau_2(y) = & F_2(y) + \int_0^y \tau_3'(z)N(x,y;1,z)dz - \int_0^1 \tau_1'(t)N(x,y,t,0)dt + \\ & + \int_0^y \left(F_2(y) + \int_0^t \tau_3'(z)N(x,y;1,z)dz - \int_0^1 \int_0^1 \tau_1'(p)N(x,y,p,0)dp \right) K(t,y)dt \end{aligned} \quad (25)$$

And finally, we put (25) into equation 2 of (21) and $\tau_2(y)$, we form the Volterra integral equation of the second type:

$$\begin{aligned} \tau_3'(y) - \int_0^y \tau_3'(z)N(1,y;1,z)dz + \int_0^y \tau_1'(z)N(1,y;t,0)dz - \int_0^z N(1,y;0,z)dz - \int_0^y \tau_3'(z)N(x,y;1,z) + \\ + \int_0^y \tau_3'(z)N(x,y;1,z)dz - \int_0^1 \int_0^1 \tau_1'(p)N(x,y,p,0)dpK(t,y)dp + \widetilde{F}_3(y) = F_2(y) \end{aligned} \quad (26)$$

There

$$\widetilde{F}_3(y) = \int_0^y F_2(z)N(1,y,0,z)dz + \int_0^y F_2(s)K(s,y)ds + \int_0^y N(1,y,0,z) \int_0^y F_2(p)K(p,y)dp \quad (27)$$

(26) tenglamani yechib $\tau_2(y)$, bundan va (25) dan $\tau_1(y)$ va (26), (27) dan $v_1(y)$ va $v_2(y)$ ni topamiz. $\tau_i(y), v_i(y) (i = \overline{1,3})$ lar ma'lum funksiyalar. Endi Ω_0 sohada I masalaning we can restore the solution, $\overline{\Omega}_i (i = \overline{1,3})$ in fields, the solution is found through the D'alambert's formula, which is the solution of the Cauchy problem. So, I-problem is solved one-valued.

The theorem is proved.



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